

# Advanced Inverse Techniques for the Design of Directional Solidification processes

Nicholas Zabaras  
Rajiv Sampath  
George Z. Yang

Sibley School of Mechanical  
and Aerospace Engineering,  
Cornell University

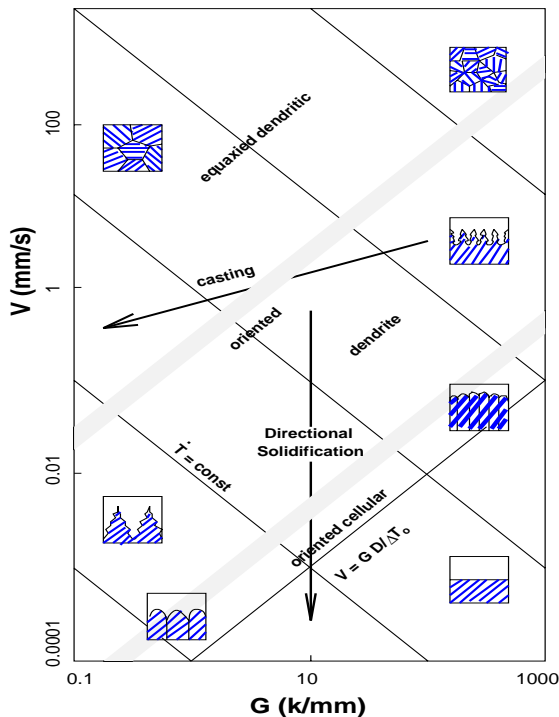
Modeling of Casting, Welding and  
Advanced Solidification Processes  
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# Outline of the presentation

- Motivation and objectives
- Background and related work
- Inverse design techniques
- Solidification design for stable growth
- Numerical results
- Forthcoming developments

# Inverse solidification design problems



Motivation:

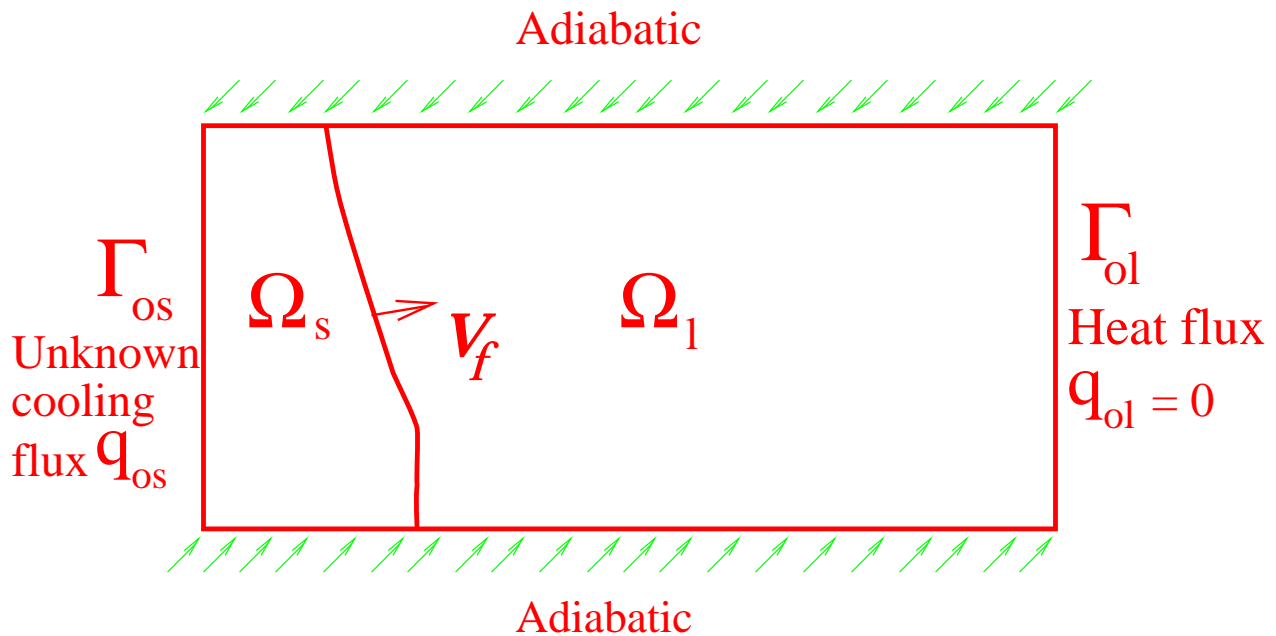
Solidification micro-  
-structures depend on  
the parameters  $G$  and  $V$

- Kurz and Fisher (1984)

Progress :

- Inverse and design Stefan problems  
Rubinski and Katz (1984)  
Zabaras (1990)  
Zabaras and Ruan (1992)
- Inverse conduction based solidification problem with data supplied from a direct melt convection problem  
Zabaras and Nguyen (1995)
- Inverse convection based solidification design problems  
Zabaras and Yang (1997, 1998)

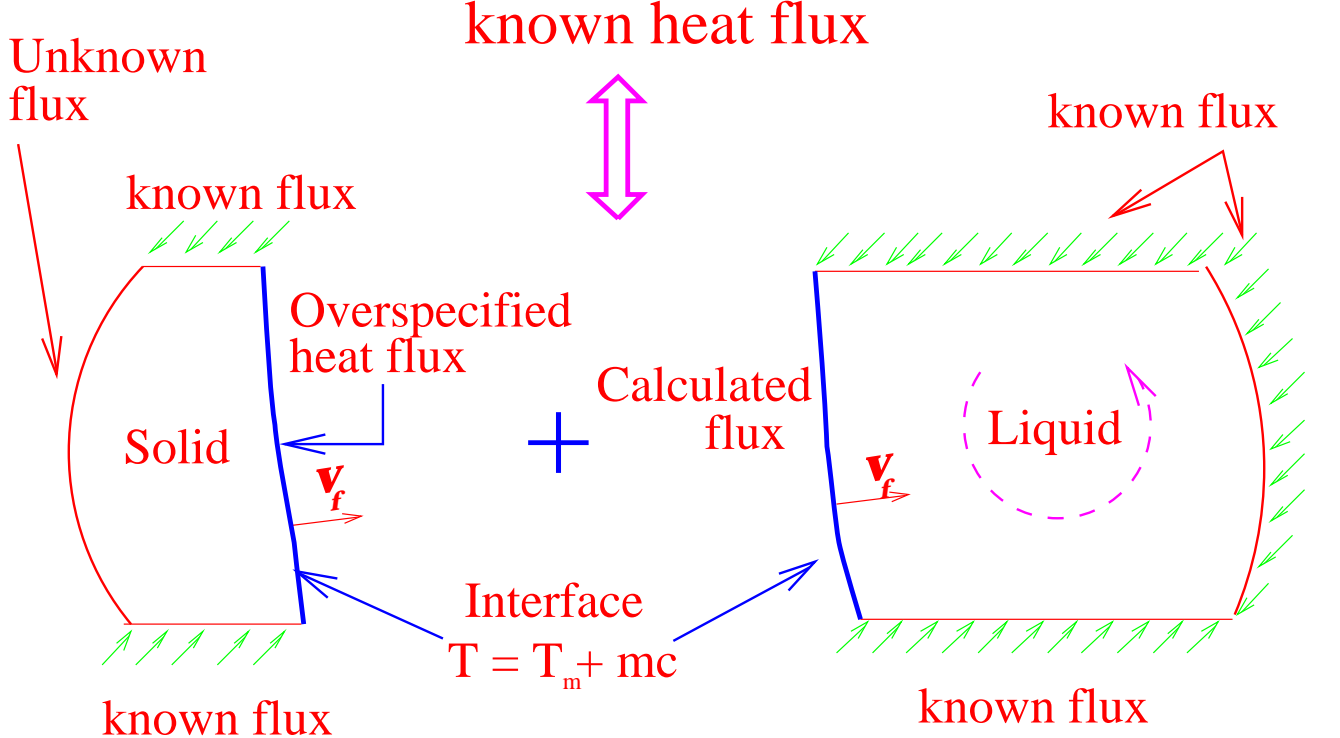
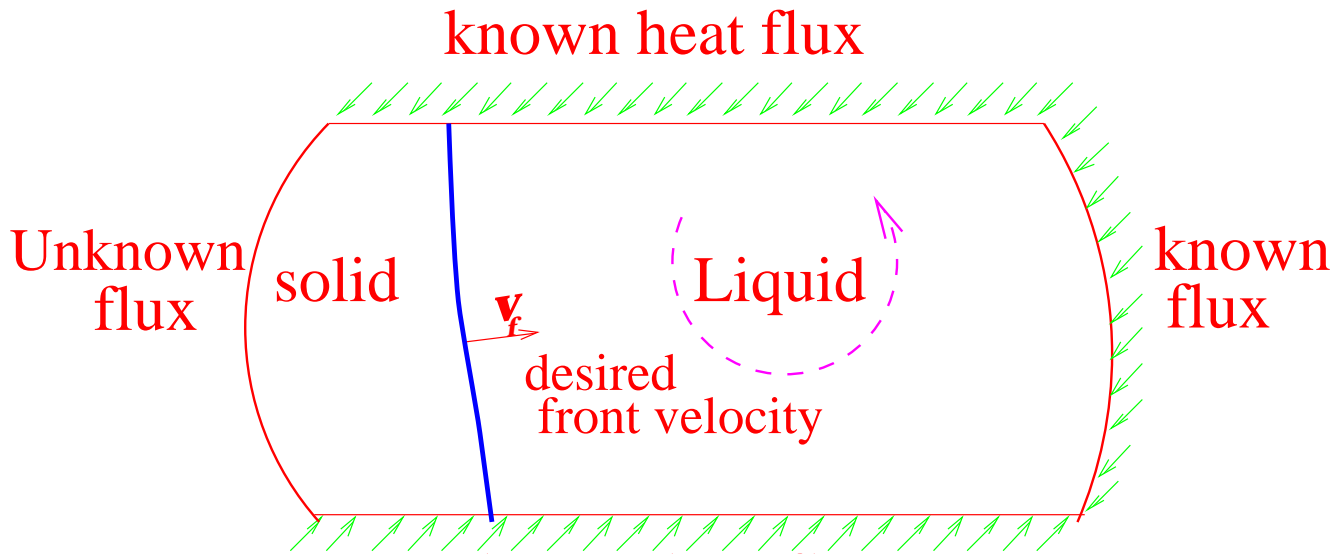
# A reference design problem in the directional solidification of a dilute binary alloy



Define the following objective:

" Find appropriate cooling flux  $q_{os}$  at the solid mold wall such that the solid-liquid interface advances with a given spatially uniform velocity  $v_f(t)$  into the liquid "

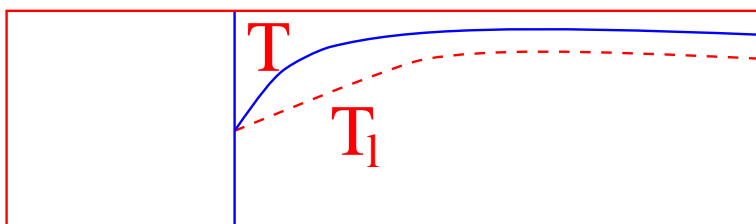
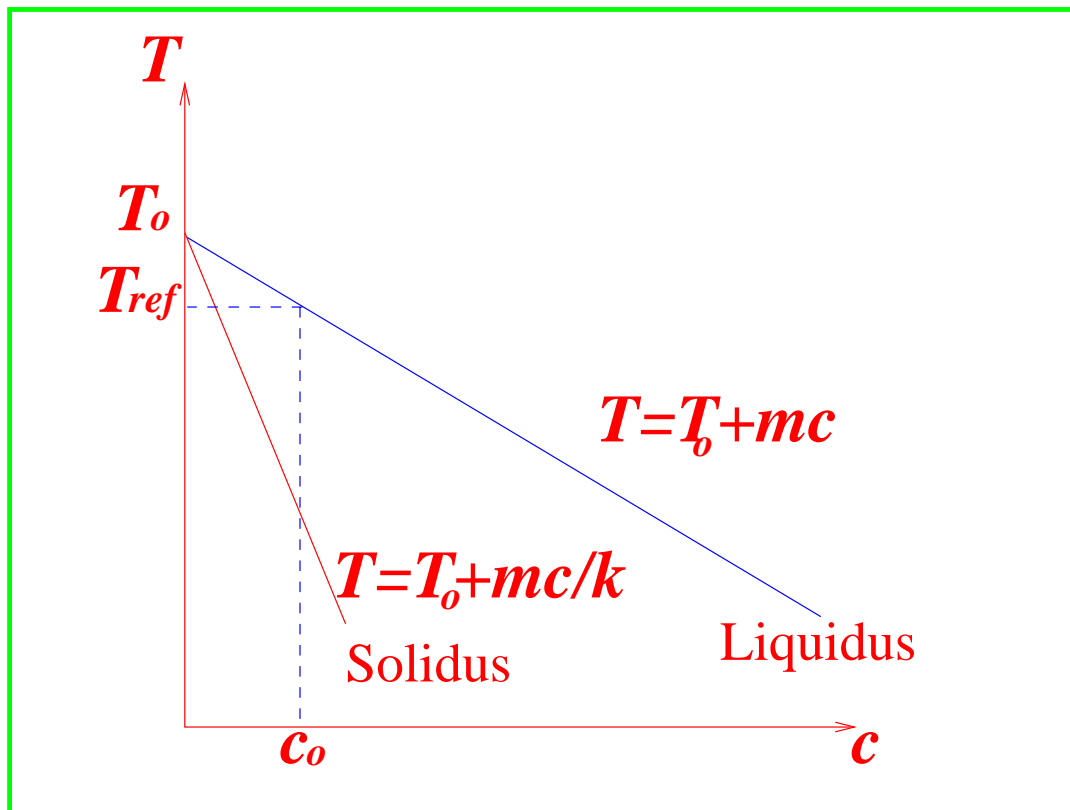
# Solution technique for the inverse reference design problem



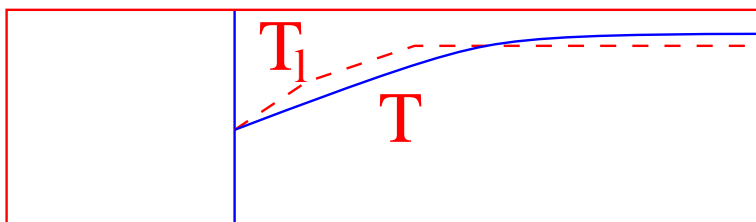
Inverse problem  
in solid

Quasi direct  
problem in liquid

# Phase diagram and constitutional stability



← Stable Growth



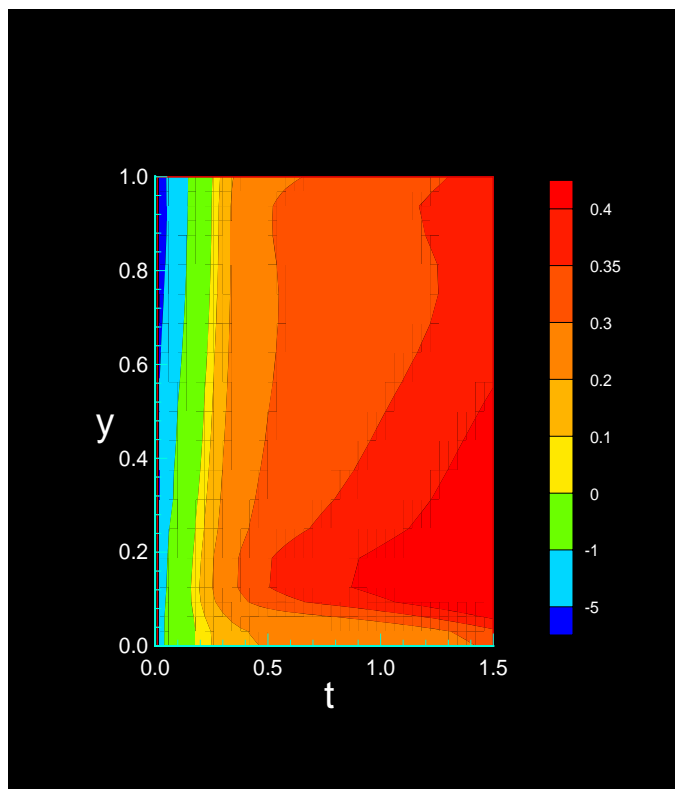
← Unstable Growth

Kurz and Fisher (1984)

# Violation of the a-priori assumption of stable growth!!

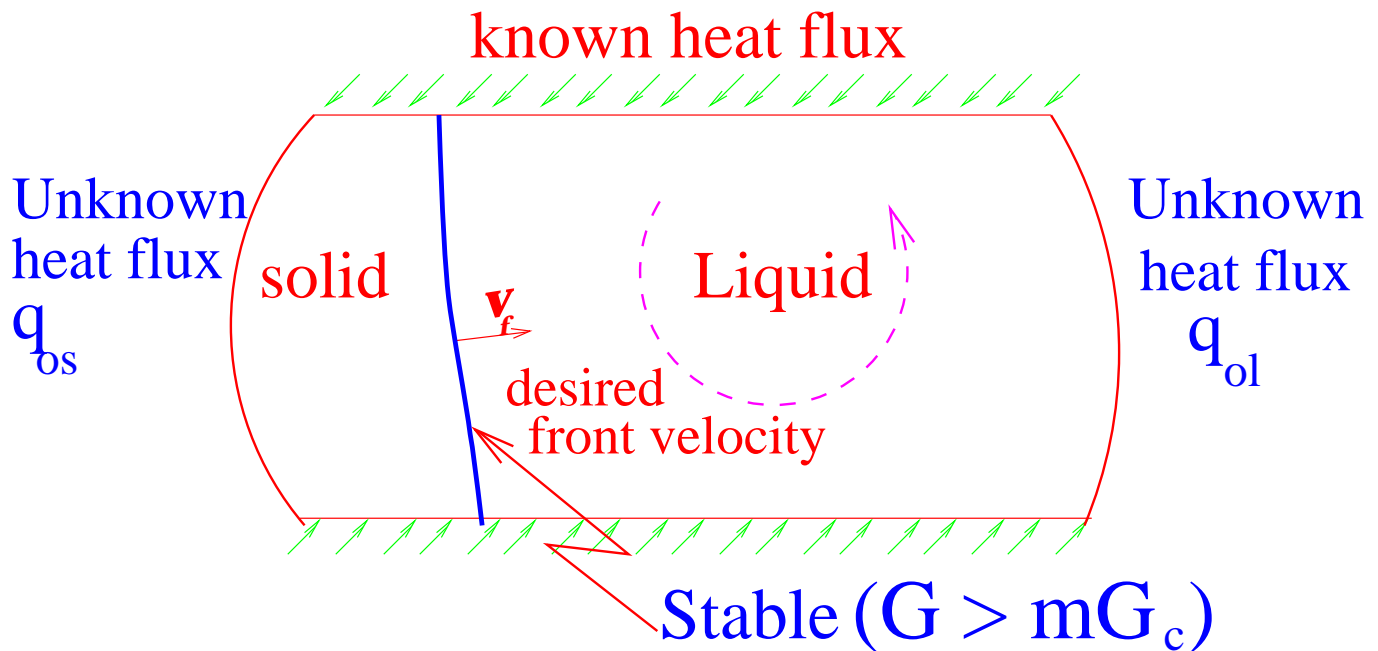
The solution of the reference design  
problem violates the constitutional  
stability criterion

$$G > mG_c \quad - \text{Kurz and Fisher (1984)}$$



Only a small region to the left is stable

# Inverse binary alloy solidification under stable desired growth conditions



Recast the inverse design problem as:

"Find the 'cooling conditions' at  $\Gamma_{os}$  *as well as* the 'heating conditions' at  $\Gamma_{ol}$  such that a *desired growth* is achieved that is ensured to be *constitutionally stable*"

## Enforcement of the stability constraint

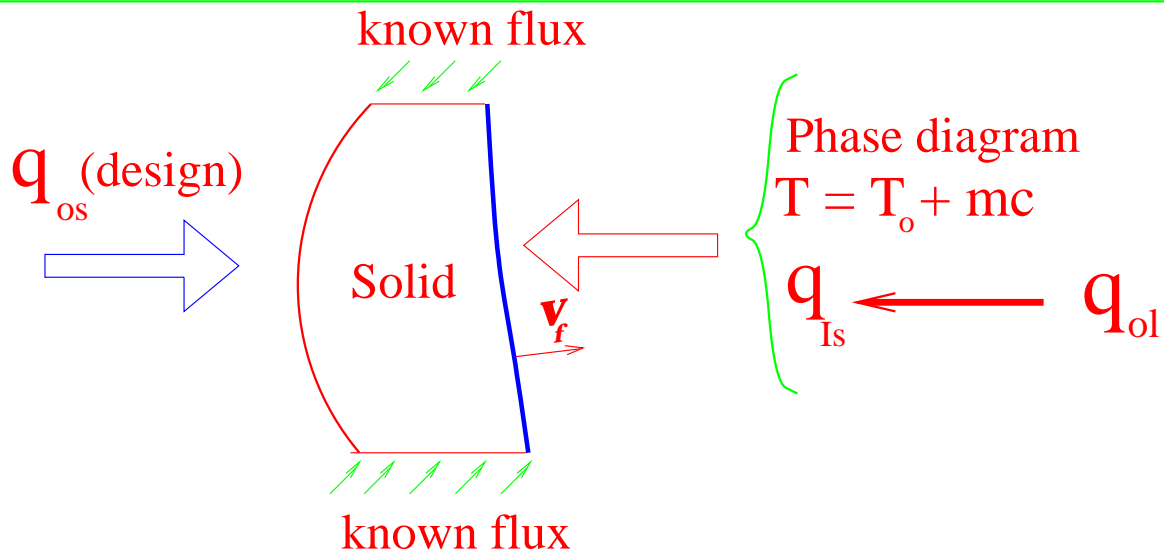
- Stability condition is enforced as follows:

$$G > mG_c \quad \longrightarrow \quad G = mG_c + \varepsilon(y,t)$$

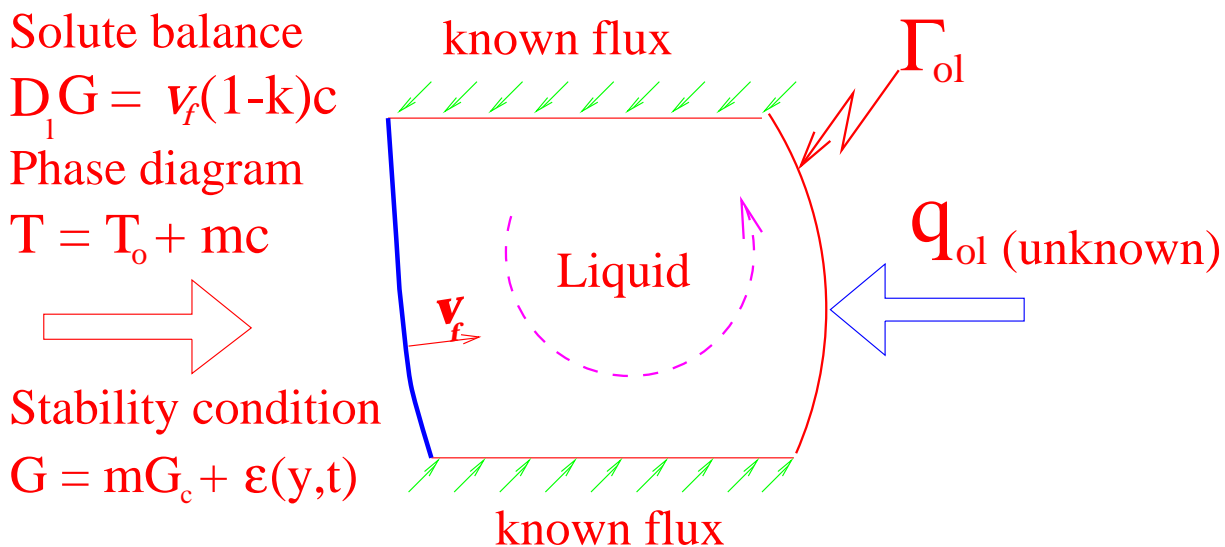
- The "over-stability" parameter  $\varepsilon(y,t)$  is here chosen so as to 'minimize' the energy input to the system

$$\varepsilon(y,t) = -m \delta(1-k) v_f c_o / D_1, \quad \delta \rightarrow 0$$

# Decompose into two sub-inverse problems



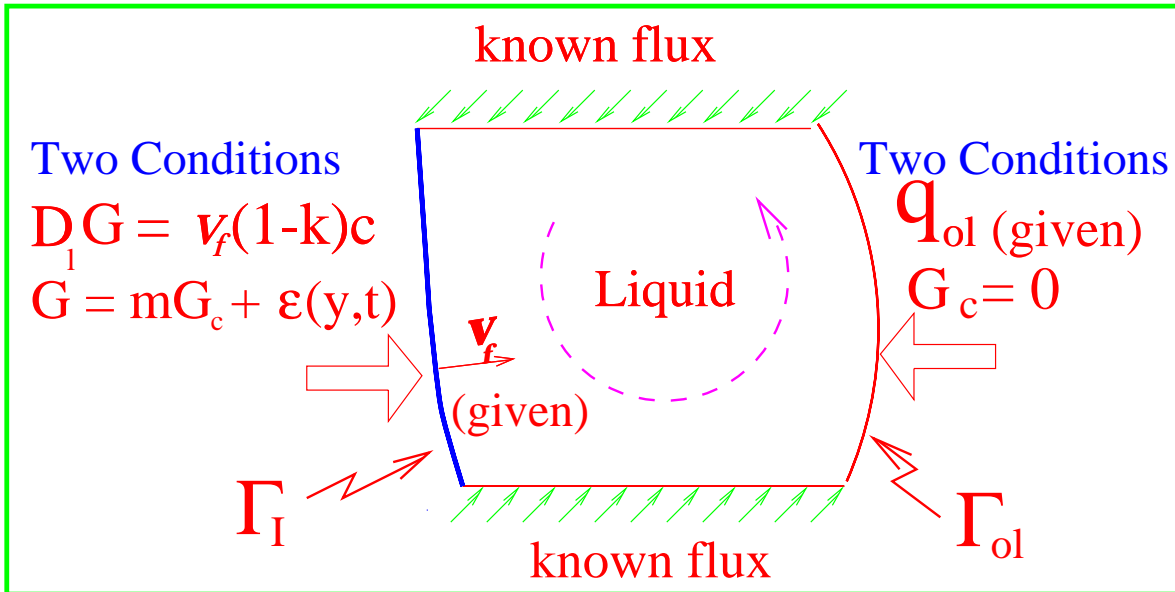
***Solid: Inverse heat conduction problem***



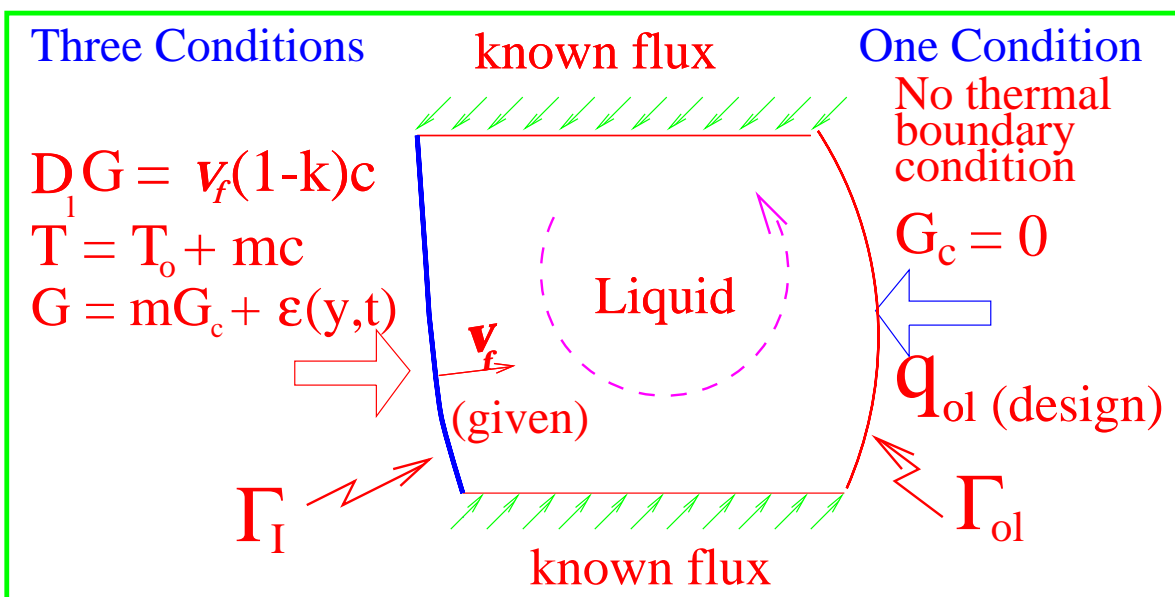
***Liquid: Inverse design with coupled heat, flow and mass transport***

# Direct versus Inverse Thermo-solutal Convection problem

*"Well-Posed Direct Problem"*

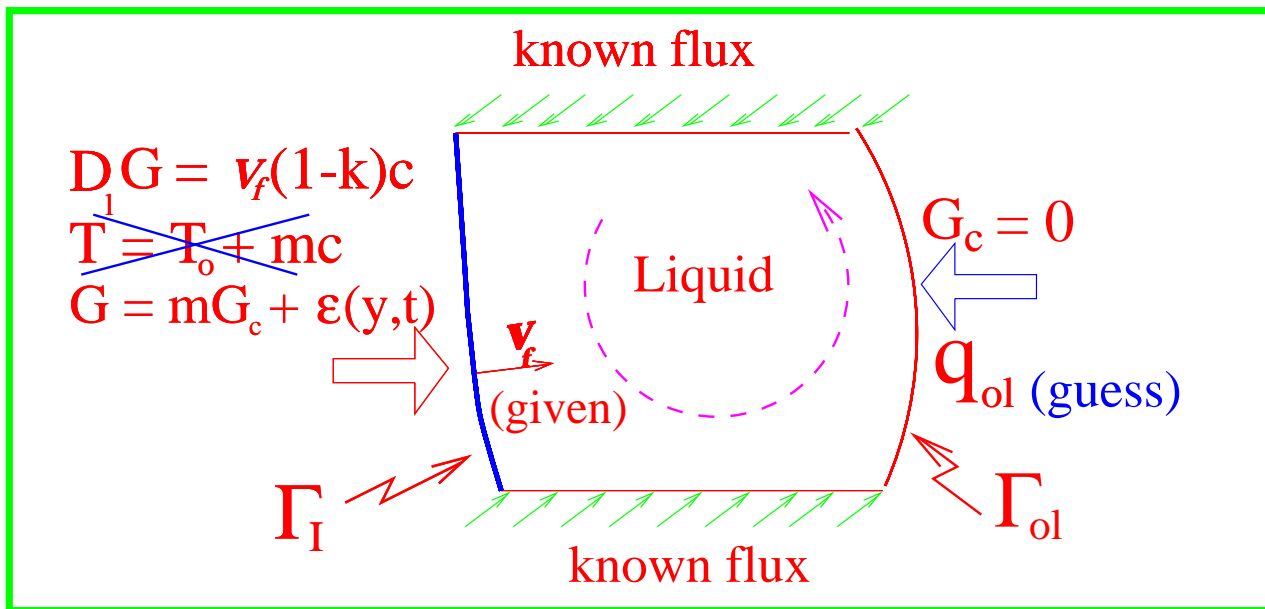


*"Ill-Posed Inverse Problem"*



# Inverse design thermo-solutal convection problem

- With a *guessed heat flux* and without using the equilibrium condition solve the following direct problem



- Define the cost functional as a measure of deviation from thermodynamic equilibrium :

$$S(q_{ol}) = \frac{1}{2} \int_0^{t_{\max}} \int_{\Gamma_I} \left[ T(\mathbf{x}, t; q_{ol}) - [T_m + mc(\mathbf{x}, t; q_{ol})] \right]^2 d\Gamma dt$$

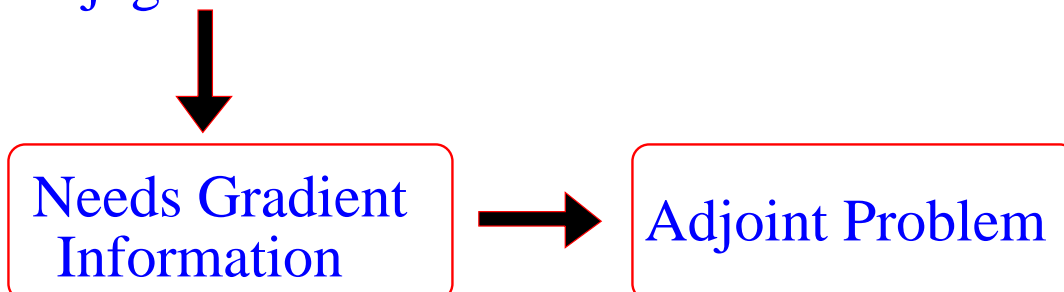
## Inverse design thermo-solutal convection problem

- Define the inverse problem in the liquid domain in a minimization sense:

Find a *quasi solution*  $\bar{q}_{ol} \in L_2(\Gamma_{ol} \times [0, t_{max}])$  such that :

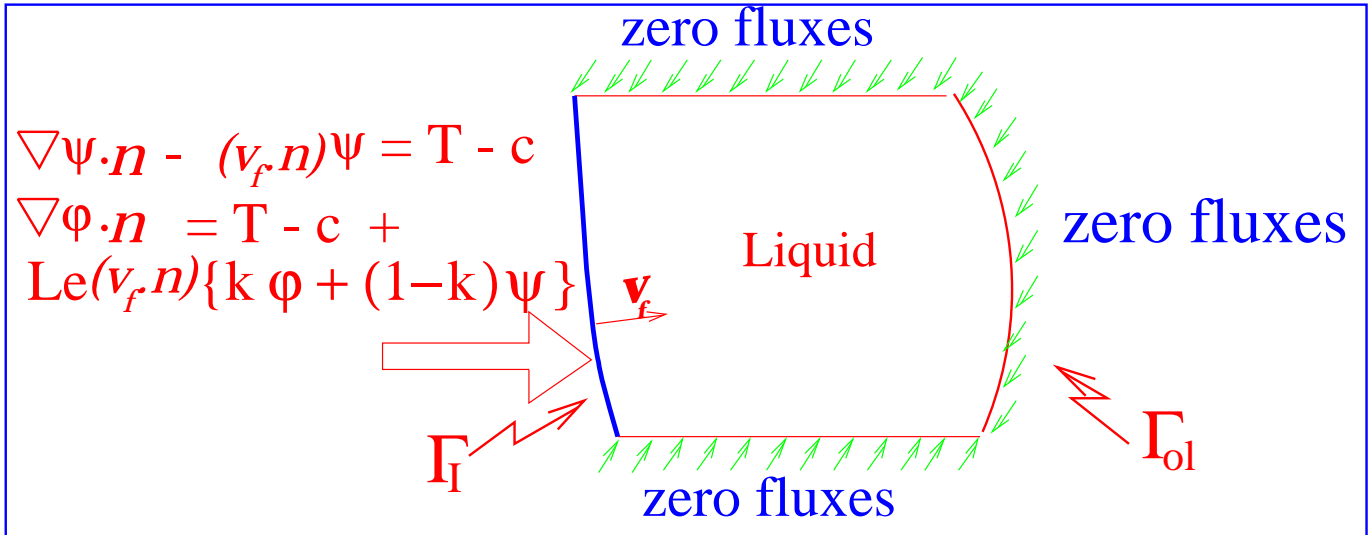
$$S(\bar{q}_{ol}) \leq S(q_{ol}) \quad \forall \quad q_{ol} \in L_2(\Gamma_{ol} \times [0, t_{max}])$$

- Solve the above minimization problem using the Conjugate Gradient method



# The adjoint problem

Gradient:  $S'(q_{ol}) = \Psi(\mathbf{x}, t; q_{ol})$  on  $\Gamma_{ol}$



Definition of adjoint fields :  $\phi(\mathbf{x}, t; q_{ol}), \psi(\mathbf{x}, t; q_{ol}), \varphi(\mathbf{x}, t; q_{ol})$

$$\frac{\partial \psi}{\partial t} + \mathbf{u} \cdot \nabla \psi = -\nabla^2 \psi + \phi \cdot \mathbf{e}_g$$

$$\frac{\partial \varphi}{\partial t} + \mathbf{u} \cdot \nabla \varphi = -\text{Le}^{-1} \left[ \nabla^2 \varphi + \frac{\gamma \text{Ra}_C}{\text{Ra}_T} \phi \cdot \mathbf{e}_g \right]$$

$$\frac{\partial \phi}{\partial t} + \mathbf{u} \cdot \nabla \phi - [\nabla \mathbf{u}]^T \phi = -\text{Pr} \nabla^2 \phi - \nabla \pi + \text{Pr} \text{Ra}_T [\psi \nabla T - \text{Le} \varphi \nabla c]$$

$$\nabla \cdot \phi = 0$$

$$\psi(\mathbf{x}, t_{max}; q_o) = \phi(\mathbf{x}, t_{max}; q_o) = \varphi(\mathbf{x}, t_{max}; q_o) = 0$$

$$\frac{\partial \psi}{\partial n} - (\mathbf{v}_f \cdot \mathbf{n}) \psi = T - c, \quad (\mathbf{x}, t) \in \Gamma_I \times [0, t_{max}]$$

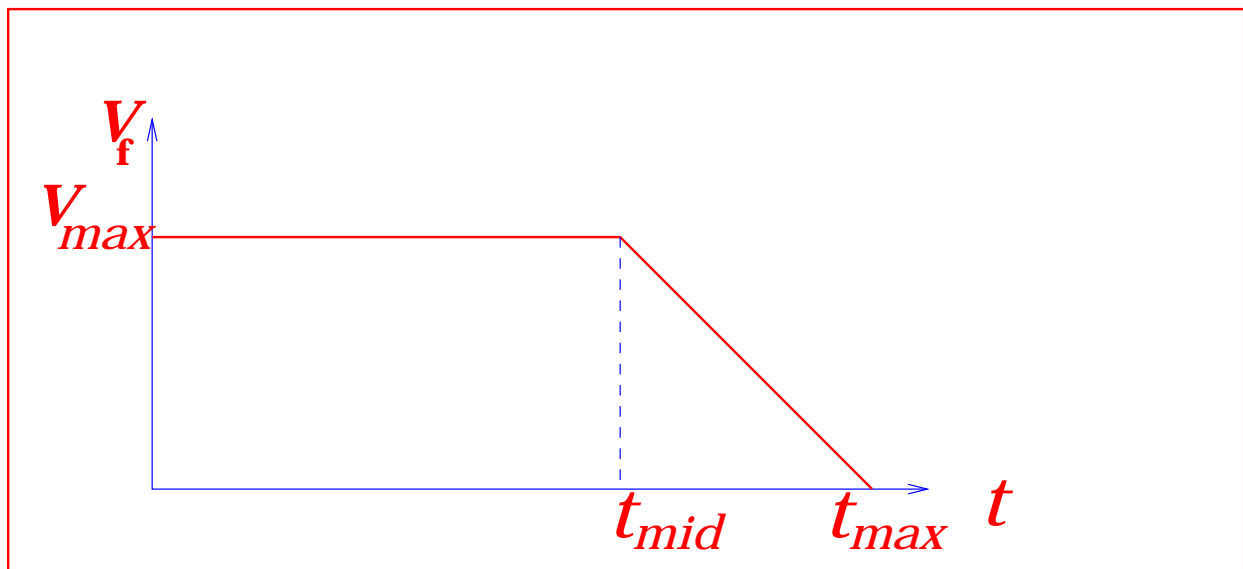
$$\frac{\partial \varphi}{\partial n} = \text{Le}(\mathbf{v}_f \cdot \mathbf{n}) [k\phi + (1-k)\psi] + T - c, \quad (\mathbf{x}, t) \in \Gamma_I \times [0, t_{max}]$$

$$\frac{\partial \psi}{\partial n} = 0, \quad \frac{\partial \varphi}{\partial n} = 0, \quad (\mathbf{x}, t) \in (\Gamma_{ol} \cup \Gamma_{hl}) \times [0, t_{max}]$$

$$\phi = 0, \quad (\mathbf{x}, t) \in \Gamma_\ell \times [0, t_{max}]$$

# Numerical Example

- Solidification of  $\text{NH}_4\text{Cl} - \text{H}_2\text{O}$  (1.5 wt% and overheat of  $20^\circ\text{C}$ )
- Design to achieve "a flat desired interface growth"
- The prescribed front velocity is chosen as shown below:



## Non-dimensionalization and parameters

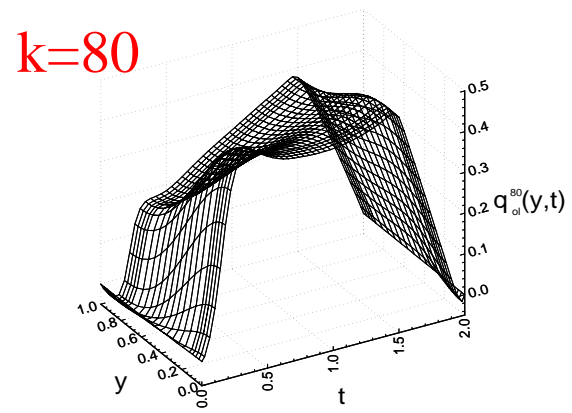
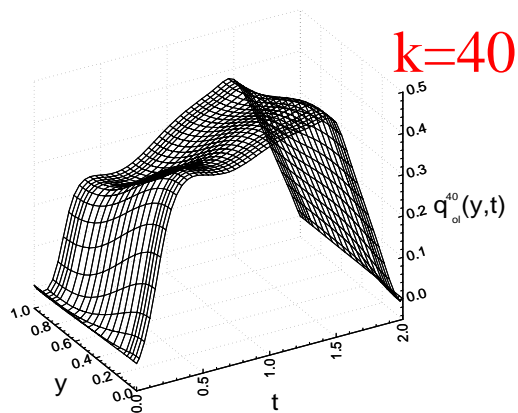
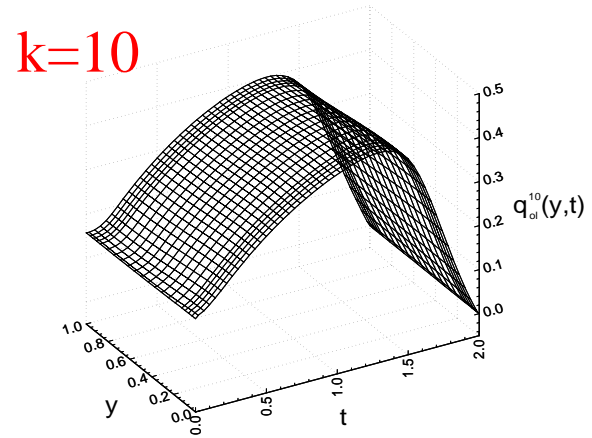
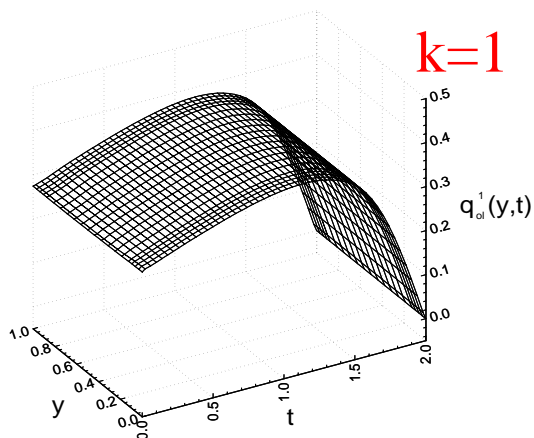
$$\bar{x} = x/L, \bar{t} = \alpha_l t/L^2, T_{\text{ref}} = T_o - mc_o$$

$$\bar{T} = \frac{T - T_{\text{ref}}}{T_i - T_{\text{ref}}}, \bar{u} = \frac{uL}{\alpha_l}, \bar{c} = \frac{c_o - c}{\gamma c_o}$$

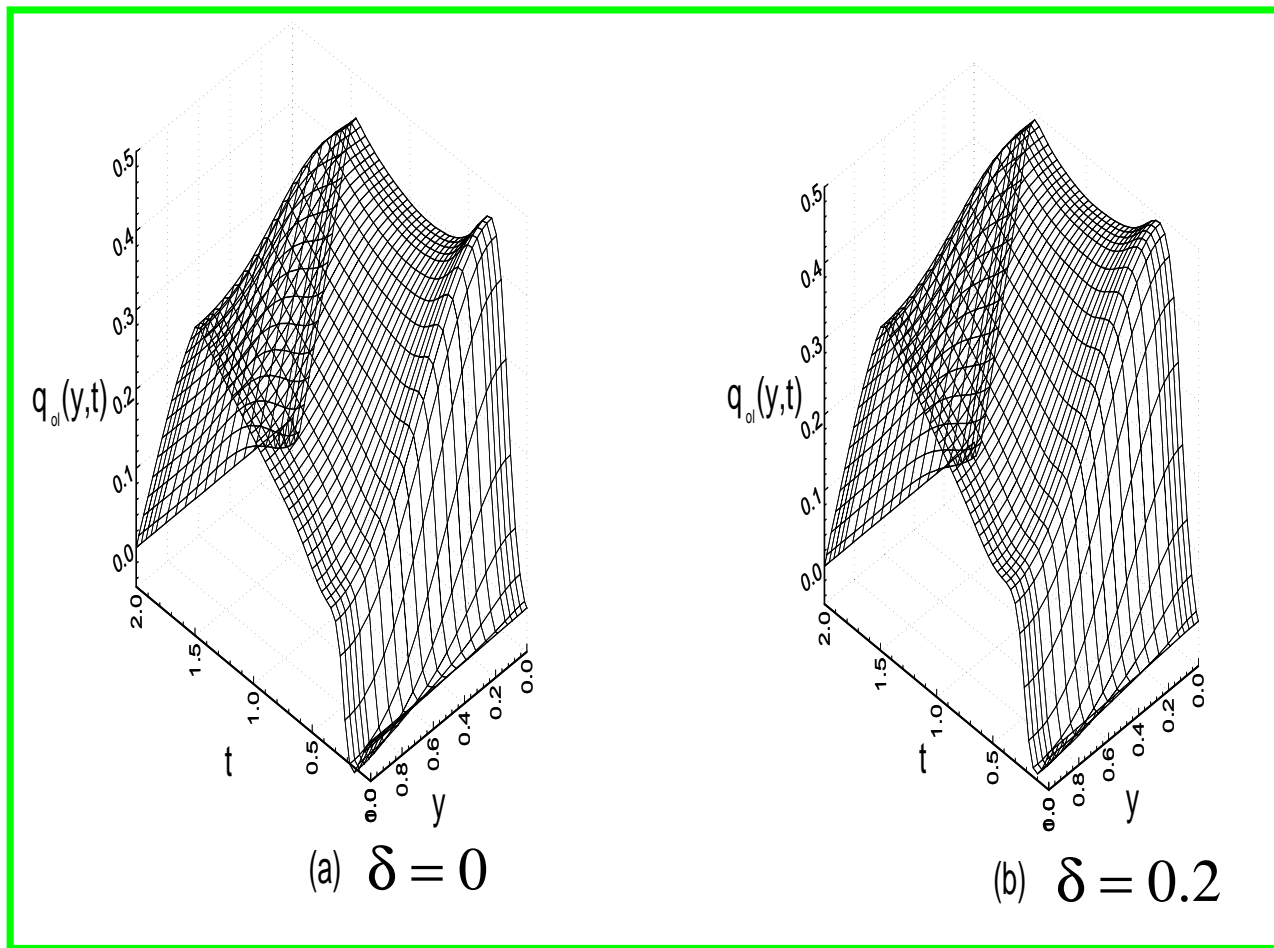
Parameter	Definition	Value
Prandtl number	$\nu/\alpha_l$	9.025
Lewis number	$\alpha_l/D_l$	27.845
Partition ratio	$c_s/c_l$	0.3
Initial overheat	$(T_i - T_o)/mc_o + 1$	18.152
Thermal Rayleigh number	$g\beta_T \Delta T L^3 / \nu \alpha_l$	$2.0 \times 10^4$
Solutal Rayleigh number	$g\beta_c \Delta c L^3 / \nu \alpha_l$	$1.0 \times 10^4$

# Heat flux solutions at intermediate iterations

\* The temporal features are captured early whereas the spatial variation needs more iterations

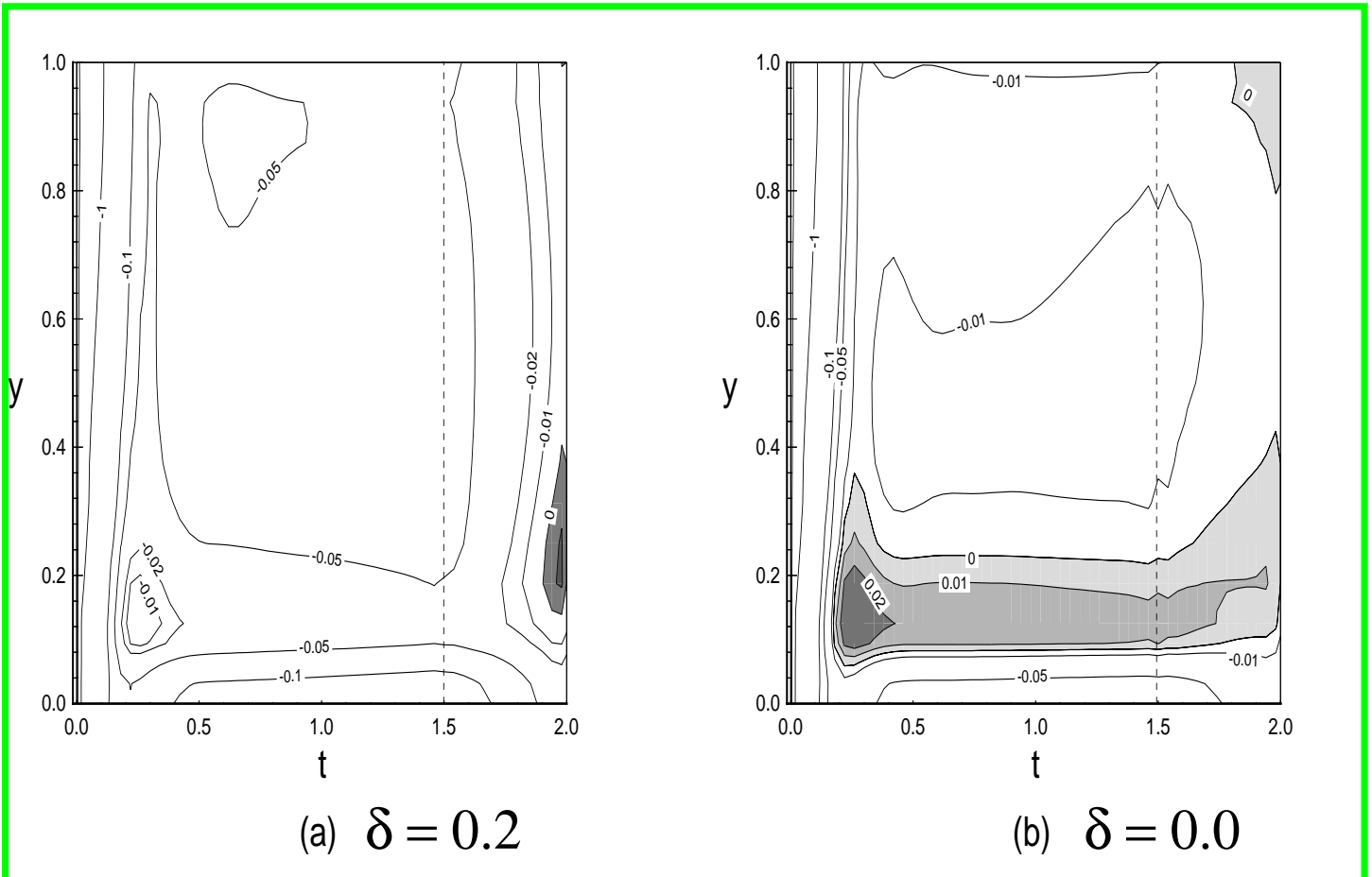


# Inverse design optimal solution: $\bar{q}_{ol}$



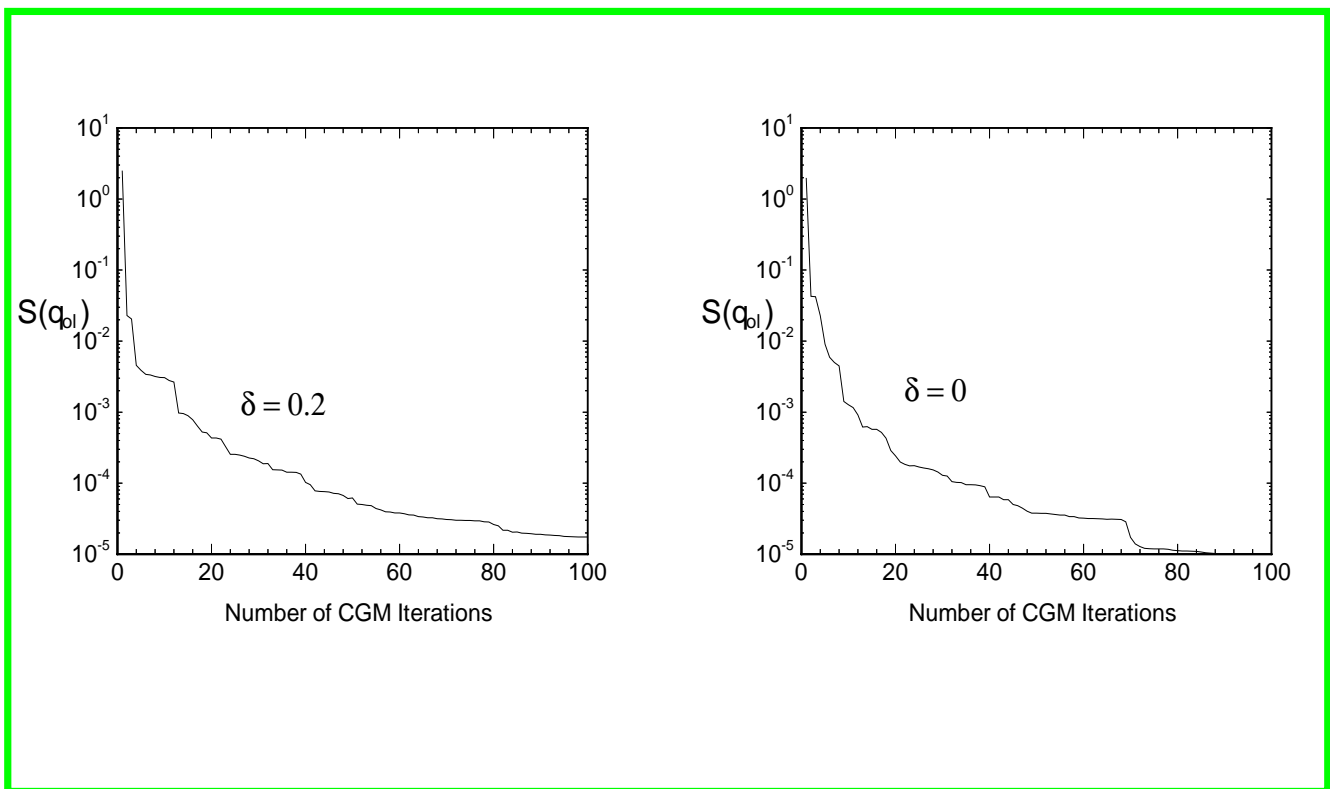
- Requires heating flux at the liquid side with greater heating at the base than at the top
- More energy input is needed as "over-stable" parameter ( $\delta$ ) increases

## Interface stability validation of the inverse solution



- \* The case  $\delta = 0$  leads to an unstable process
- \*  $\delta > 0$  is thus necessary to maintain stability during the "whole" process

# Convergence of the inverse design algorithm



- \* Rapid drop in the cost functional initially but slower at later steps
- \* Convergence is better for lower values of  $\delta$

# Determination of the solid side heat flux $q_{os}$

Additional  
Parameters

Inverse  
problem  
in the  
liquid

$q_{ol}$

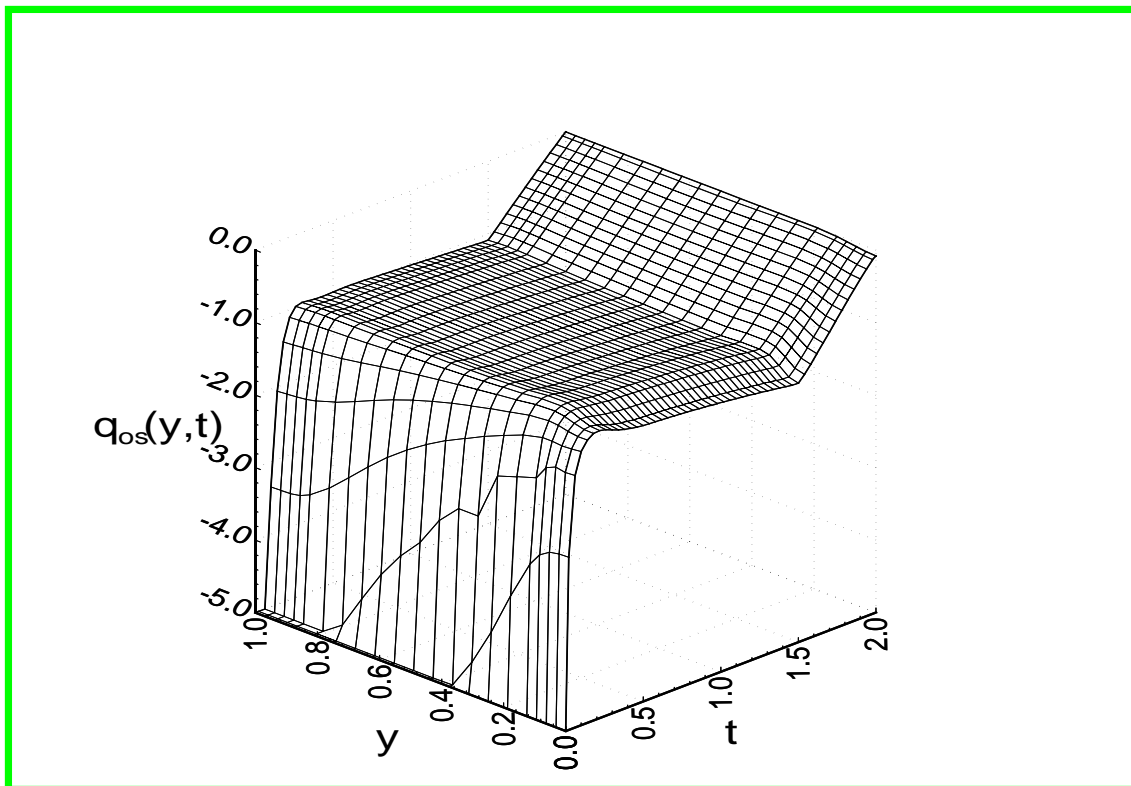
Inverse problem  
in the solid

$$R_k = 1$$

$$R_\alpha = 1$$

$$Ste = 0.3$$

$q_{os}$



# Salient features and limitations

## Salient features :

- ☆ New innovative adjoint technique to control binary alloy solidification systems
- ☆ Control parameters chosen based on factors governing the formation of microstructures
- ☆ Enforces stability to ensure physically valid results

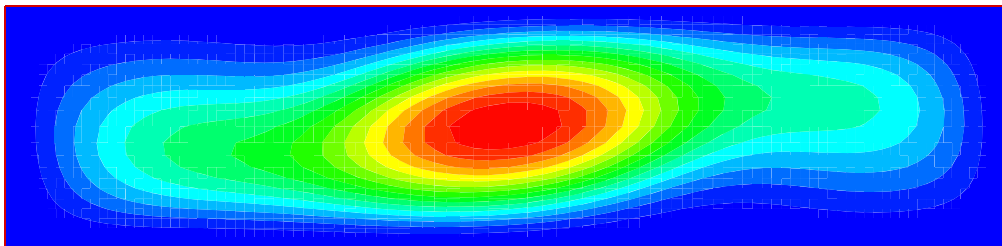
## Limitations :

- ☆ Industrial implementation of optimal heat flux solution is quite challenging
- ☆ Only a "quasi solution" is guaranteed
- ☆ The "optimal solution" may lead to side effects that are undesirable

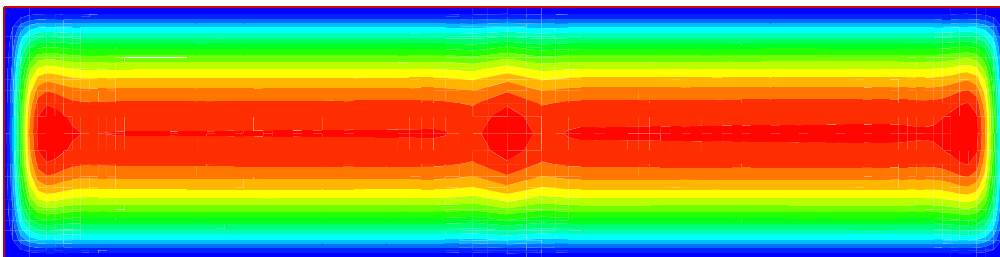
## Forthcoming developments

### Thermo-Electromagnetic control:

- ☆ The presence of a strong magnetic field effectively damps the melt flow



$$\text{Ha} = 0, \Psi_{\max} = 0.55$$

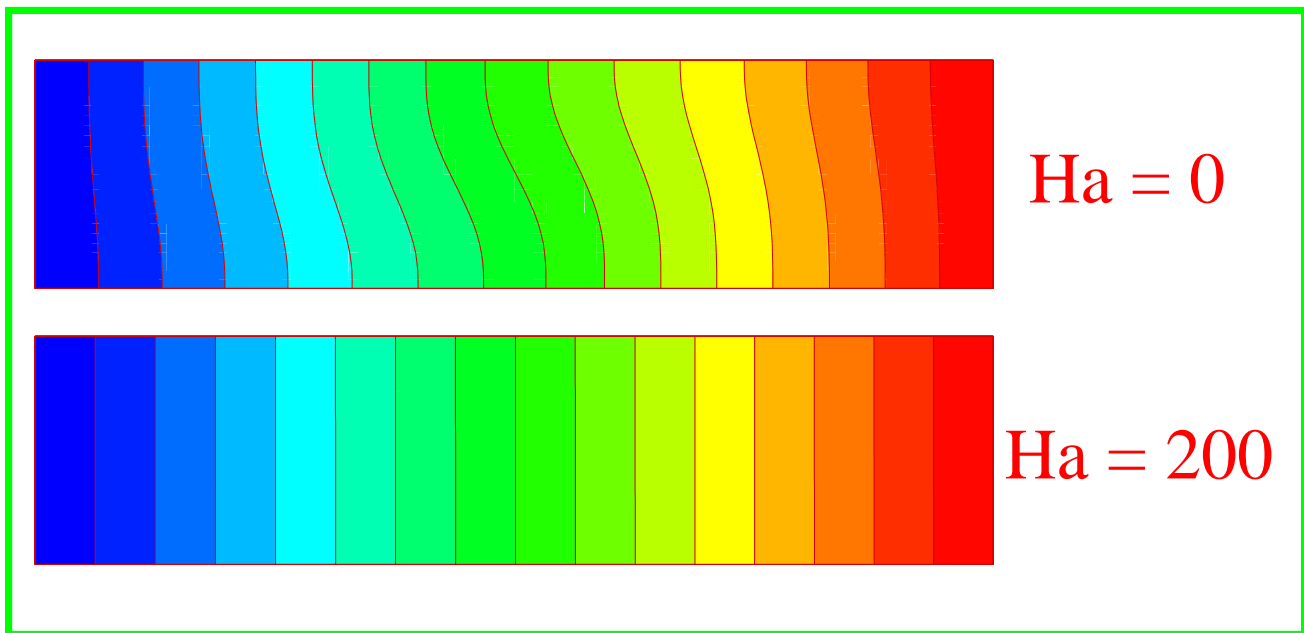


$$\text{Ha} = 200, \Psi_{\max} = 0.00063$$

Stream function plot showing distinct drop in flow circulation

## Forthcoming developments

- ☆ Magnetic control can make the optimal heat flux solutions smoother as the damping of the flow reduces the curving of the interface.



Curved isotherms get flatter with increasing Hartmann number

- ☆ Due to magnetic damping we can expect to have lesser solute variation in the crystal