

## HOMEWORK 5

### LDA of DFT for the Calculation of the Ground State of the Helium Atom

**Due: Monday March 2nd, 4:30 p.m.**

In this assignment you will need to write a code (by modifying the Hartree-Fock code in HW3) to calculate the ground state of a helium atom using the local density (LDA) approximation to the density functional theory (DFT).

The Kohn-Sham wave-function for He satisfies the equation

$$\left( -\frac{1}{2}\nabla^2 + v_{ext}(r) + v_h(r) + v_{xc}(r) \right) \phi(r) = \varepsilon \phi(r)$$

where  $v_{ext}(r) = -\frac{2}{r}$  is the external potential due to the nucleus,  $v_h(r) = \int \frac{n(r')}{|r-r'|} d^3r'$  is the Hartree potential and  $v_{xc}(r)$  is the exchange-correlation potential, which depends on the approximation. In this homework, we will use LDA (see lecture notes) in which the exchange-correlation energy is written as:

$$E_{xc}^{LDA} = \int n(r) [\varepsilon_x(n(r)) + \varepsilon_c(n(r))] d^3r$$

where the exchange energy density  $\varepsilon_x^{LDA}$  and the correlation energy density  $\varepsilon_c^{LDA}$  (Perdew-Zunger parametrization) are given [in the introduction to DFT lecture notes](#) on the course web site in terms of the Wigner-Seitz radius  $r_s$ . To obtain the potential, notice that  $v_{xc} = \frac{\delta E_{xc}^{LDA}}{\delta n} = \frac{\partial(n\varepsilon_{xc})}{\partial n}$ .

As with the Hartree-Fock method, we will use here a Gaussian basis:

$$\phi = \sum_i C_i \chi_i(r), \quad \chi_i(r) = e^{-\alpha_i r^2}$$

in which the Kohn-Sham equations are written in a matrix form as follows:

$$\sum_j (h^{ij} + V_h^{ij} + V_{xc}^{ij}) C_j = \varepsilon \sum_j S^{ij} C_j$$

where  $h^{ij}$  and  $S^{ij}$  are the same as in the Hartree-Fock implementation, and  $V_h^{ij}$  and  $V_{xc}^{ij}$  are the matrix elements of the Hartree and exchange-correlation potentials, respectively.

Modify the HF codes to implement these calculations and compute the ground state of the He atom. The evaluation of the matrix elements  $V_h^{ij}$  can be done using the analytical expression  $V_h^{ij} = 2 \sum_{kl} C_k C_l Q_{ijkl}$  which should be familiar from your HF program or using

numerical integration (quadrature) to approximate  $V_h^{ij} = \int \chi_i^*(r) v_h(r) \chi_j(r) d^3r$ , where the Coulomb potential is the solution of the Poisson equation  $\nabla^2 v_h(r) = -4\pi n(r)$ . Note that the cost of the first approximation is  $\sim O(N^4)$ , whereas the second scales as  $\sim O(N^2)$ , where  $N$  is the number of the Gaussian basis functions.

To compute  $V_{xc}^{ij}$ , transform the integral of the form  $V_{xc}^{ij} = \int_0^\infty 4\pi r^2 F(r) dr$  to a domain where

you can do numerical integration, e.g. to  $V_{xc}^{ij} = \int_{-1}^{+1} f(x) \sqrt{1-x^2} dx$ , using  $r = \frac{\sqrt{1+x}}{\sqrt{1-x}}$ . For

this case, you can perform the integration in  $x$  using the [Gauss-Chebyshev quadrature of the 2<sup>nd</sup> kind](#) (see [A First Course in Numerical Analysis](#), R. Rabinowitz, pp. 111):

$$\int_{-1}^{+1} f(x) \sqrt{1-x^2} dx = \sum_{j=1}^n f(a_j) H_j, \text{ where } H_j = \frac{\pi}{n+1} \sin^2 \frac{j\pi}{n+1}, \text{ and}$$

$$a_j = \cos \frac{j\pi}{n+1}, j = 1, 2, \dots, n.$$